

Problème inverse dans une équation de Schrödinger

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▶ Introduction

- ▶ Usual Schrödinger equation
- ▶ Non-linear inverse problem
- ▶ Uniqueness and Stability
- ▶ Transmission Schrödinger equation

▶ Method to prove stability for some linear inverse problem

▶ Carleman estimates

- ▶ Weight function
- ▶ Proof scheme of the usual case
- ▶ Transmission case

▶ Stability Theorem

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Schrödinger equation

$T > 0$, $\Omega \subset \mathbb{R}^2$ bounded domain

Smooth boundary $\partial\Omega$

potential $q = q(x) \in L^\infty(\Omega)$

$$y = y(x, t)$$

$$y' = \frac{\partial y}{\partial t}$$

$$\begin{cases} iy'(x, t) + \operatorname{div}(a(x)\nabla y(x, t)) + q(x)y(x, t) = 0, & x \in \Omega, t \in (0, T) \\ y(x, t) = h(x, t), & x \in \partial\Omega, t \in (0, T) \\ y(x, 0) = y_0(x), & x \in \Omega. \end{cases}$$

If $a = 1$, $y_0 \in L^2(\Omega)$, $h \in L^2(\partial\Omega \times (0, T))$, $\exists!$ solution

$$y = y(q) \in C([0, T], H^{-1}(\Omega)) \cap H^{-1}(0, T; L^2(\Omega))$$

such that $\frac{\partial y}{\partial \nu} \in H^{-2}([0, T], H^{-\frac{3}{2}}(\partial\Omega))$.

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Nonlinear inverse problem :

$$\begin{cases} iy' + \operatorname{div}(a(x)\nabla y) + qy = 0, & \Omega \times (0, T) \\ y = h, & \partial\Omega \times (0, T) \\ y(0) = y_0, & \Omega. \end{cases}$$

Is it possible to retrieve the potential $q = q(x)$, $x \in \Omega$ from measurement of the normal derivative $\frac{\partial y}{\partial \nu} \Big|_{\partial\Omega \times (0, T)}$?

- ▶ $q \mapsto \frac{\partial y(q)}{\partial \nu}$ is non linear
- ▶ Local result (q known, p unknown, close to q)
- ▶ This is a one-measurement inverse problem.

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Uniqueness and stability

- ▶ Uniqueness :

$$\left(\frac{\partial y(p)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)} = \frac{\partial y(q)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)} \right) \Rightarrow (p = q \text{ sur } \Omega) ?$$

- ▶ Stability : It is possible to estimate $(q - p)|_{\Omega}$ by

$$\frac{\partial y(q)}{\partial \nu} - \frac{\partial y(p)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)}$$

in suitable norms ?

Schrödinger transmission equation

Ω_1 , smooth boundary Γ_1 st $\bar{\Omega}_1 \subset \Omega$ and $\Omega_0 = \Omega \setminus \bar{\Omega}_1$.

We set $a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases}$ with $a_j > 0$ for $j = 0, 1$

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Inverse problem well-posed provided that :

- ▶ Ω_1 is strongly convex
ie Ω_1 is convex and Γ_1 has a non-vanishing curvature
- ▶ $a_0 < a_1$

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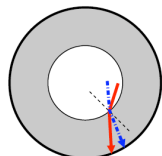
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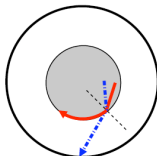
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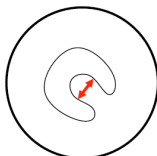
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$a_1 > a_0$ (good)



$a_1 < a_0$ (bad)



Ω_1 non convex (bad)

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References

- ▶ Inverse problem with Carleman estimates
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- ▶ Wave equation, constant main coefficient
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Method

Takes its roots in the works of Bukhgeim, Klibanov and Malinski.

- ▶ Equation in $y(q)$
- ▶ Equation in $z = u'$
- ▶ Equation in $u = y(q) - y(p)$
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$$\begin{cases} iy' - \operatorname{div}(a\nabla y) + qy = 0, & \Omega \times (0, T) \\ y = h, & \partial\Omega \times (0, T) \\ y(0) = y_0, & \Omega \end{cases} \quad \begin{cases} iu' - \operatorname{div}(a\nabla u) + (f + q)u = fR \\ u = 0 \\ u(0) = 0. \end{cases}$$

$$\begin{cases} u = y(q) - y(p) \\ f = p - q \\ R = y(q) \end{cases} \Rightarrow \text{Linearized version of the inverse problem :}$$

Can we determine f in Ω from $\frac{\partial u}{\partial \nu} \Big|_{\Omega \times (0, T)}$?

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Extension of the equation on $(-T, T)$

- ▶ Since $R(0) \in \mathbb{R}$, we extend v and R on $\Omega \times (-T, 0)$ by $v(x, t) = -\bar{v}(x, -t)$ and $R(x, t) = \bar{R}(x, -t)$.
 $z(x, t) = -\bar{z}(x, -t)$, $R(x, t) = \bar{R}(x, -t)$ on $(-T, 0)$
- ▶ If we assume $R(0) \in i\mathbb{R}$, an other extension can be defined.

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- ▶ If we assume $R(0) \in i\mathbb{R}$, an other extension can be defined.

We take $q \in L^\infty(\Omega)$ and $u_0 \in H_0^1(\Omega)$ such that

- ▶ $R = y(q) \in H^1(0, T; L^\infty(\Omega))$
- ▶ $0 < r < |R(x, 0)|$ a.e. on Ω

and we will use an estimate like

$$C_1(s) \int_{\Omega} |z'(0)|^2 dx \leq C_2(s) \int_{-T}^T \int_{\partial\Omega} \left| \frac{\partial z}{\partial \nu} \right|^2 d\sigma dt + C_3(s) \int_{-T}^T \int_{\Omega} |\text{rhs}|^2 dx$$

with $\frac{C_3(s)}{C_1(s)} \rightarrow 0$ when $s \rightarrow \infty$.

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Carleman Estimate for the Schrödinger equation

$Lv = iv' + \Delta v + qv$ we suppose that $\|q\|_{L^\infty(\Omega)} \leq m$.

For λ and s large enough, $\exists M = M(\Omega, T, m) > 0$ such that :

$$\begin{aligned} & s\lambda \int_{-T}^T \int_{\Omega} |\nabla v|^2 \varphi e^{-2s\varphi} dxdt + s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |v|^2 \varphi^3 e^{-2s\varphi} dxdt \\ & \leq M \int_{-T}^T \int_{\Omega} |Lv|^2 \varphi^3 e^{-2s\varphi} dxdt + Ms\lambda \int_{-T}^T \int_{\partial\Omega} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \end{aligned}$$

$\forall v \in L^2(-T, T; H_0^1(\Omega))$ such that $Lv \in L^2(\Omega \times (-T, T))$ and $\frac{\partial v}{\partial \nu} \in L^2(\Omega \times (-T, T))$.

Schrödinger weight function

$$\varphi(x, t) = \frac{\alpha - e^{\lambda\psi(x)}}{(T-t)(T+t)}$$

- ▶ $\lambda > 0$, $\|e^{\lambda\psi}\|_{L^\infty} \leq \alpha$
- ▶ $\psi > 0$ regular on Ω
- ▶ $|\nabla\psi| \geq \beta > 0$ in Ω
- ▶ $\exists \varepsilon > 0$ such that $\forall \xi \in \mathbb{R}^N$, $\lambda|\nabla\psi \cdot \xi|^2 + D^2\psi(\xi, \bar{\xi}) \geq \varepsilon|\xi|^2$
- ▶ $\nabla\psi \cdot \nu < 0$ on $\partial\Omega \setminus \Gamma_0$

ex : $\psi(x) = |x - x_0|^2$, $x_0 \in \mathbb{R}^N \setminus \bar{\Omega}$

We set $w = e^{-s\varphi}v$ and we calculate $Pw = e^{-s\varphi}L(e^{s\varphi}w)$.

We get $Pw = P_1w + P_2w + qw$ with

$$\begin{cases} P_1w = iw' + \Delta w + s^2|\nabla\varphi|^2w \\ P_2w = is\varphi'w + 2s\nabla\varphi \cdot \nabla w + s\Delta\varphi w. \end{cases}$$

$$\begin{aligned} & \int_{-T}^T \int_{\Omega} |Pw - qw|^2 dxdt \\ &= \int_{-T}^T \int_{\Omega} (|P_1w|^2 + |P_2w|^2) dxdt + 2 \operatorname{Re} \int_{-T}^T \int_{\Omega} P_1w \overline{P_2w} dxdt. \end{aligned}$$

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The point is to obtain in the calculus :

- ▶ some “dominant” terms with the right sign

$$s\lambda \int_{-T}^T \int_{\Omega} |\nabla w|^2 dxdt, \quad s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |w|^2 dxdt$$

- ▶ and the “measurement” term

$$s\lambda \int_{-T}^T \int_{\Gamma_0} \left| \frac{\partial w}{\partial \nu} \right|^2 \nabla \psi \cdot \nu d\sigma dt$$

Discontinuous main coefficient

$$a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases} \quad \text{with } a_j > 0 \text{ for } j = 0, 1.$$

Then for each $f \in L^2(\Omega \times (0, T))$, y satisfies the equation

$$iy' - \operatorname{div}(a\nabla y) + qy = f \quad \text{in } \Omega \times (0, T)$$

iff $j \in \{0, 1\}$, $y_j = y|_{\Omega_j}$ satisfies

$$iy'_j - a_j \Delta y_j + qy_j = f|_{Q_j} \quad \text{in } \Omega_j \times (0, T)$$

with the *transmission conditions* on the interface $\Gamma_1 = \partial\Omega_1$

$$\begin{cases} y_0 = y_1 \\ a_0 \frac{\partial y_0}{\partial \nu_0} + a_1 \frac{\partial y_1}{\partial \nu_1} = 0 \end{cases} \quad \text{on } \Gamma_1 \times (0, T).$$

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Carleman estimate - discontinuous a

We use a weight-function $\varphi(x, t) = \frac{\alpha - e^{\lambda\psi(x)}}{(T-t)(T+t)}$ such that :

- ▶ **Condition 1** : $\varphi|_{\Omega_j} =$ appropriate weight-function in Ω_j .
- ▶ **Condition 2** : We can control the “interface” terms on $\partial\Omega_1$.
- ▶ The choice of ψ is crucial :

$$\begin{cases} |\nabla\psi| \geq \beta > 0 \\ 2D_a^2\psi(\xi, \bar{\xi}) + 2a^2\lambda|\nabla\psi \cdot \xi|^2 - a\nabla a \cdot \nabla\psi|\xi|^2 \geq \epsilon|\xi|^2, \quad \forall \xi \in \mathbb{C}^n \end{cases} \text{ in } \Omega_0 \cup \Omega_1$$

where $D_a^2\psi \left(a \frac{\partial}{\partial x_i} \left(a \frac{\partial \psi}{\partial x_j} \right) \right)_{1 \leq i, j \leq N}$

$$\begin{cases} \psi_0 = \psi_1 = \text{cst} \\ a_0 \frac{\partial \psi_0}{\partial \nu_0} + a_1 \frac{\partial \psi_1}{\partial \nu_1} = 0 \text{ and } \frac{\partial \psi_0}{\partial \nu_0} + \frac{\partial \psi_1}{\partial \nu_1} < 0 \end{cases} \text{ on } \Gamma_1$$

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$$\begin{cases} |\nabla\psi| \geq \beta > 0 \\ 2D_a^2\psi(\xi, \bar{\xi}) + 2a^2\lambda|\nabla\psi \cdot \xi|^2 - a\nabla a \cdot \nabla\psi|\xi|^2 \geq \epsilon|\xi|^2, \quad \forall \xi \in \mathbb{C}^n \end{cases} \quad \text{in } \Omega_0 \cup \Omega_1$$

where $D_a^2\psi \left(a \frac{\partial}{\partial x_i} \left(a \frac{\partial \psi}{\partial x_j} \right) \right)_{1 \leq i, j \leq N}$

$$\begin{cases} \psi_0 = \psi_1 = cst \\ a_0 \frac{\partial \psi_0}{\partial \nu_0} + a_1 \frac{\partial \psi_1}{\partial \nu_1} = 0 \quad \text{and} \quad \frac{\partial \psi_0}{\partial \nu_0} + \frac{\partial \psi_1}{\partial \nu_1} < 0 \end{cases} \quad \text{on } \Gamma_1$$

Carleman estimate - discontinuous a

We use a weight-function $\varphi(x, t) = \frac{\alpha - e^{\lambda\psi(x)}}{(T-t)(T+t)}$ such that :

- ▶ **Condition 1** : $\varphi|_{\Omega_j} =$ appropriate weight-function in Ω_j .
- ▶ **Condition 2** : We can control the “interface” terms on $\partial\Omega_1$.
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$$\begin{cases} |\nabla\psi| \geq \beta > 0 \\ 2D_a^2\psi(\xi, \bar{\xi}) + 2a^2\lambda|\nabla\psi \cdot \xi|^2 - a\nabla a \cdot \nabla\psi|\xi|^2 \geq \epsilon|\xi|^2, \quad \forall \xi \in \mathbb{C}^n \end{cases} \quad \text{in } \Omega_0 \cup \Omega_1$$

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Explicit weight function

We will actually define an ε -pair of transmission weight functions $(\psi^1, \psi^2) \in C^4(\Omega_0 \cup \Omega_1)^2$ such that for each $j, k \in \{1, 2\}$ with $j \neq k$,

$$\psi^j - \psi^k \geq \delta > 0 \quad \text{in } B_\varepsilon(x_k)$$

and

$$\left\{ \begin{array}{l} |\nabla \psi^j| \geq \beta > 0 \\ \forall \xi \in \mathbb{C}^n, \\ 2D_a^2 \psi^j(\xi, \bar{\xi}) + 2a^2 \lambda |\nabla \psi^j \cdot \xi|^2 - a \nabla a \cdot \nabla \psi^j |\xi|^2 \geq \epsilon |\xi|^2 \end{array} \right. \quad \text{in } \Omega_0 \cup \Omega_1 \setminus B_\varepsilon(x_j)$$

$$\psi^j(x) = \eta(x) \frac{a_k}{\rho(x)^2} |x - x_j|^2 + M_j, \quad x_j \in \Omega_1$$

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Intermediate step

$$Z = \left\{ v \in L^2(-T, T; H_0^1(\Omega)) : Lv \in L^2(Q), \frac{\partial v}{\partial \nu} \in L^2(\Sigma), v \text{ satisfies } (Tr) \right\},$$

$$\|w\|_{s,\lambda,\psi} = s^3 \lambda^4 \int_{-T}^T \int_{\Omega} \theta^3 |w|^2 dx dt + s \lambda \int_{-T}^T \int_{\Omega} \theta |\nabla w|^2 dx dt$$

and $\|\cdot\|_{s,\lambda,\psi,B}$ corresponds to the above terms defined in $B \subset \Omega$.

We prove first that

$$\begin{aligned} & \|P_1(w)\|_{L^2}^2 + \|P_2(w)\|_{L^2}^2 + \|w\|_{s,\lambda,\psi}^2 \\ & \leq C \|P(w)\|_{L^2}^2 + C \|w\|_{s,\lambda,\psi,B_\varepsilon(x_j)}^2 + s \lambda C \iint_{\partial\Omega} \left| a \frac{\partial w}{\partial \nu} \right|^2. \end{aligned}$$

Carleman estimate

Suppose there exists for some $\varepsilon > 0$ an ε -pair of transmission weight functions (ψ^1, ψ^2) belonging to $C^4(\Omega_0 \cup \Omega_1)$.

$\forall q \in L^\infty(\Omega)$ with $\|q\|_{L^\infty(\Omega)} \leq m$, $\exists C = C(\Omega, T, m) > 0$ such that

$$\begin{aligned} & \sum_{k=1}^2 \left(s\lambda \int_{-T}^T \int_{\Omega} |\nabla v|^2 e^{-2s\varphi^k} dx dt + s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |v|^2 e^{-2s\varphi^k} dx dt \right) \\ & \leq C \sum_{k=1}^2 \left(\int_{-T}^T \int_{\Omega} |Lv|^2 e^{-2s\varphi^k} dx dt + s\lambda \int_0^T \int_{\partial\Omega} \left| a_1 \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \right) \end{aligned}$$

$\forall v \in L^2(-T, T; H_0^1(\Omega))$ satisfying the transmission conditions... and for all λ and s large enough.

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$\forall v \in L^2(-T, T; H_0^1(\Omega))$ satisfying the transmission conditions... and for all λ and s large enough.

Stability theorem

$$\begin{cases} iy' + \operatorname{div}(a(x)\nabla y) + qy = 0 \\ y = h \\ y(0) = y_0 \end{cases}$$

Let \mathcal{U} be a bounded subset of $L^\infty(\Omega)$, $q \in L^\infty(\Omega)$, $T_0 > 0$ and $r > 0$.

If $y_0 \in H_0^1(\Omega)$ and y_0 takes its values in \mathbb{R} and if

$$\begin{aligned} |y_0(x)| &\geq r, \text{ a.e. in } \Omega, \\ y(q) &\in H^1(0, T; L^\infty(\Omega)), \end{aligned}$$

then $\exists C = C(\Omega_0, \Omega_1, T, a_0, a_1, \|q\|_{L^\infty}, y_0, h, \mathcal{U}, r) > 0$ such that $\forall p \in \mathcal{U}$, we have

$$\|q - p\|_{L^2(\Omega)} \leq C \left\| a_1 \frac{\partial y(q)}{\partial \nu} - a_1 \frac{\partial y(p)}{\partial \nu} \right\|_{H^1(0, T; L^2(\Gamma_0))}.$$

$$\begin{cases} iz' - \operatorname{div}(a\nabla z) + qz = fR' \\ z = 0 \\ z(0) = -ifR(0) \end{cases}$$

We set $w = e^{s\varphi} z$ and define

$$I = \operatorname{Im} \int_{-T}^0 \int_{\Omega} P_1 w \bar{w} \, dx dt.$$

where $P_1 w = iw' + \Delta w + s^2 |\nabla \varphi|^2 w$.

► Calculation :

$$\begin{aligned} I &= \operatorname{Im} \int_{-T}^0 \int_{\Omega} (iw' + \Delta w + s^2 |\nabla \varphi|^2 w) \bar{w} \, dx dt \\ &= \frac{1}{2} \int_{\Omega} |w(x, 0)|^2 \, dx \\ &= \frac{1}{2} \int_{\Omega} |R(x, 0)|^2 e^{-2s\varphi(0)} |f|^2 \, dx \end{aligned}$$

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► Carleman :

$$\begin{aligned}
 I &\leq \left(\int_{-T}^T \int_{\Omega} |P_1 w|^2 dx dt \right)^{\frac{1}{2}} \left(\int_{-T}^T \int_{\Omega} |w|^2 dx dt \right)^{\frac{1}{2}} \\
 &\leq M s^{-\frac{3}{2}} \left(\int_0^T \int_{\Omega} |P w|^2 dx dt + s \int_0^T \int_{\Gamma_0} \left| \frac{\partial w}{\partial \nu} \right|^2 d\sigma dt \right) \\
 &\leq M s^{-\frac{3}{2}} \left(\int_0^T \int_{\Omega} |f R'|^2 e^{-2s\varphi} dx dt + s \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \right)
 \end{aligned}$$

Since $|R(x, 0)| \geq r > 0$, then

$\exists g_0 \in L^2(0, T)$, $|R'(x, t)| \leq g_0(t)|R(x, 0)|$ so that

$$\begin{aligned}
 \int_{\Omega} |f|^2 |R(0)|^2 e^{-2s\varphi(0)} dx &\leq M s^{-\frac{3}{2}} \int_0^T \int_{\Omega} |f|^2 |g_0|^2 |R(0)|^2 e^{-2s\varphi(0)} dx dt \\
 &+ M s^{-\frac{1}{2}} \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt.
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$$\begin{aligned} I &\leq \left(\int_{-T}^T \int_{\Omega} |P_1 w|^2 dx dt \right)^{\frac{1}{2}} \left(\int_{-T}^T \int_{\Omega} |w|^2 dx dt \right)^{\frac{1}{2}} \\ &\leq M s^{-\frac{3}{2}} \left(\int_0^T \int_{\Omega} |P w|^2 dx dt + s \int_0^T \int_{\Gamma_0} \left| \frac{\partial w}{\partial \nu} \right|^2 d\sigma dt \right) \\ &\leq M s^{-\frac{3}{2}} \left(\int_0^T \int_{\Omega} |f R'|^2 e^{-2s\varphi} dx dt + s \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \right) \end{aligned}$$

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We finally obtain

$$(1 - Cs^{-\frac{3}{2}}) \int_{\Omega} |f(x)|^2 dx \leq C \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt,$$

and for s large enough, we prove the stability of the linear inverse problem :

$$\|f\|_{L^2(\Omega)} \leq C \left\| \left\| \frac{\partial u}{\partial \nu} \right\| \right\|_{H^1(0,T;L^2(\Gamma_0))}.$$

Theorem - Linear inverse problem

Let $\|q\|_{L^\infty} \leq m$ and $r > 0$. We assume that $a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases}$ with $a_1 > a_0 > 0$ and

$$\begin{aligned} R(0) &\in \mathbb{R} \\ |R(x, 0)| &\geq r \quad \text{a.e. in } \Omega, \\ R &\in H^1(0, T; L^\infty(\Omega)). \end{aligned}$$

Then $\exists C = C(\Omega_0, \Omega_1, T, \Gamma_0, m, \|R\|) > 0$ such that $\forall f \in L^2(\Omega)$, the solution u of

$$\begin{cases} iu' + \operatorname{div}(a\nabla u) + qu = fR, & \Omega \times (0, T) \\ u = 0, & \partial\Omega \times (0, T) \\ u(0) = 0, & \Omega, \end{cases}$$

satisfies

$$\|f\|_{L^2(\Omega)} \leq C \left\| a_1 \frac{\partial u}{\partial \nu} \right\|_{H^1(0, T; L^2(\partial\Omega))}.$$

Further problems

- ▶ Recovering the main coefficient.
- ▶ Observations in a part of the boundary.
- ▶ Logarithmic stability (without geometric hypothesis).
- ▶ Neumann boundary data and Dirichlet measurement.

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